

Electrostatic energy in dielectric Media

Let, the charge q is on the positive plate, so that the potential difference is q/c . The total work, then to go from $q=0$ to $q=Q$ is,

$$W = \int_0^Q \frac{q}{c} dq = \frac{1}{2} CV^2 \quad (\text{electrostatic energy in vacuum})$$

where, $Q = CV$.

If the capacitor is filled with linear dielectric, its capacitance exceeds the vacuum value by a factor of the dielectric constant,

$$C = \epsilon_r C_{vac}$$

Hence, work will be increased by same factor. The reason is that you have to pump on more ~~free~~ free charge to achieve a given potential, bcoz part of the field is cancelled by the bound charges. (Hence, for a dielectric filled capacitor,

$$W = \frac{\epsilon_0}{2} \int \epsilon_r \cdot E^2 dz = \frac{1}{2} \epsilon_0 \epsilon_r \int |E|^2 dz \quad (1)$$

Since, energy stored in electrostatic system is

$$W = \frac{\epsilon_0}{2} \int |E|^2 dz \quad (2)$$

then eqn (1) reduces to;

$$W = \frac{1}{2} \int \underline{D} \cdot \underline{E} dz$$

(in presence of linear dielectrics)

To prove it, suppose the dielectric material is fixed in position and we bring in the free charge q , bit at a time. As P_f is increased by an amount ΔP_f , the polarization will change; (The work done on the incremental free charge:

$$\Delta W = \int (\Delta P_f) V dz$$

Since, $\nabla \cdot \underline{D} = P_f$, $\Delta P_f = \nabla \cdot (\Delta \underline{D})$ where $\Delta \underline{D}$ is the resulting change in \underline{D} , so

$$\Delta W = \int [\nabla \cdot (\Delta \underline{D})] V dz$$

Now,

$$\nabla \cdot [(\Delta \underline{D}) V] = [\nabla \cdot (\Delta \underline{D})] V + \Delta \underline{D} \cdot (\nabla V)$$

and hence,

$$\Delta W = \int_V \nabla \cdot [(\Delta \underline{D}) V] dz + \int_V (\Delta \underline{D}) \cdot (\nabla V) dz$$

$$\because \underline{E} = -\nabla V$$

$$\Delta W = \int_V \nabla \cdot [(\Delta \underline{D}) V] dz + \int_V (\Delta \underline{D}) \cdot \underline{E} dz$$

using div. Theorem

$$\int_V \nabla \cdot \underline{A} dz = \oint_S \underline{A} \cdot \hat{n} ds$$

Then,

$$\Delta W = \oint_S (\Delta \underline{D} \cdot V) \cdot \hat{n} ds + \int_V (\Delta \underline{D}) \cdot \underline{E} dz$$

then,

$$\Delta W = \int_V (\Delta \underline{D}) \cdot \underline{E} dz \quad \text{--- (3)}$$

surface integral will be vanish

Now, If the medium is a linear dielectric, then $\underline{D} = \epsilon \underline{E}$, So

$$\frac{1}{2} \Delta (\underline{D} \cdot \underline{E}) = \frac{1}{2} \Delta (\epsilon E^2) = \epsilon (\Delta \underline{E}) \cdot \underline{E} \quad \cancel{\epsilon \cdot \underline{E} \cdot \Delta \underline{E}}$$

$$\text{Thus, by (3) or, } \frac{1}{2} \Delta (\underline{D} \cdot \underline{E}) = \underline{(\Delta \underline{D}) \cdot \underline{E}}$$

$$\Delta W = \Delta \left(\frac{1}{2} \int \underline{D} \cdot \underline{E} dz \right)$$

Then, the total work done,

$$\Delta W = \frac{1}{2} \int \underline{D} \cdot \underline{E} dz \quad \text{--- (4)}$$

ϵ_0 (2) is apply in case of vacuum and ϵ_0 (4) is apply in presence of dielectric. ϵ_0 (2) does not include the work done in stretching and "spring" the dielectric molecules it does not include the spring energy $\frac{1}{2}kx^2$. With the unpolarized dielectric in place, we bring in the free charges one by one allowing the dielectric to respond as it sees fit. Since free charge is actually push around then we get ϵ_0 (4).

The total energy of the system

$$W_{\text{total}} = W_{\text{free}} + W_{\text{bound}} + W_{\text{spring}}$$

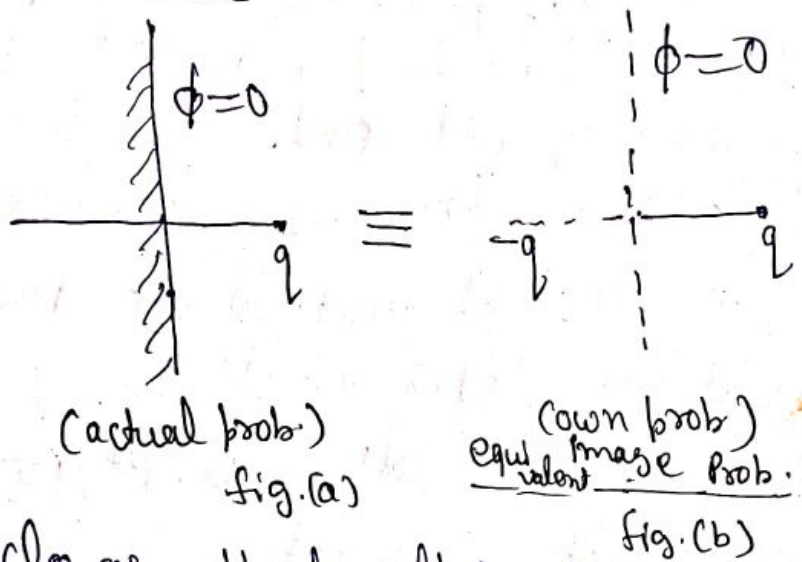
In Method (2) last two are equal and opposite, thus Method (2) in calculating W_{free} , actually delivers W_{tot} , whereas Method (1), by calculating $W_{\text{free}} + W_{\text{bound}}$, leaves out W_{spring} .

Method of Images: - ^{the method of images} concerns itself with one or more point charges in the presence of boundary surfaces. ex:- Conductors either grounded or held at fixed potential. Under Condⁿ it is possible that suitably charges of appropriate magnitude external to region of interest can produce boundary Condⁿ.

These charges are called image charges. and replacement of actual prob with boundaries by an enlarged region with image charges is called Method of images. Image charges must be external to the volume

of interests so ~~there~~ their potential must be Sol^n inside the volume. Particular integral is Sol^n of Poisson eqs. is provided by ^{Poiss} sum of potential of charges inside the volume.

Ex: - A point charge is located in front of this ∞ conductor at zero potential



It is clear that this is equivalent ^{original} to the prob. of actual charge and ~~equivalent~~ equal and opposite charge ^{located at} mirror image point ~~define~~ ^{in the} position of conductor behind the plane defined by the position of the conductor.